

Boundary conditions

	Heat eq ⁿ	Wave eq ⁿ	Math name
$\Phi = 0$	Fixed temperature	Fixed displacement	Dirichlet BC or BC of the 1st kind
$\frac{d\Phi}{dx} = 0$	Insulated	Free end	Neumann BC or BC of the 2nd kind
$\frac{d\Phi}{dx} = \pm h\Phi$	Newton's law of Cooling	Elastic.	Mixed BC or Robin BC or BC of the 3rd kind
$\Phi(-L) = \Phi(L)$ and $\frac{d\Phi(-L)}{dx} = \frac{d\Phi(L)}{dx}$	periodic	periodic	periodic
$ \Phi(0) < \infty$	bounded at the origin	bounded at the origin	singularity condition

Thm Regular Sturm-Liouville problem

Def Regular Sturm-Liouville problem is

$$\frac{d}{dx} \left(p(x) \frac{d\phi}{dx} \right) + q(x) \phi + \lambda \sigma(x) \phi = 0, \quad a < x < b$$

$$\left. \begin{aligned} \beta_1 \phi(a) + \beta_2 \frac{d\phi}{dx}(a) &= 0 \\ \beta_3 \phi(b) + \beta_4 \frac{d\phi}{dx}(b) &= 0 \end{aligned} \right\}$$

General BCs:

1st, 2nd and 3rd kinds of BCs.

β_1, \dots, β_4 : real numbers

Functions $p(x)$, $q(x)$, $\sigma(x)$ are real-valued and

continuous on $a \leq x \leq b$. In addition, $p(x) > 0$, $q(x) > 0$
and $a \leq x \leq b$.

Thm Regular Sturm-Liouville problem

1. All eigenvalues λ are real.
2. There are infinitely many eigenvalues
 $\lambda_1 < \lambda_2 < \lambda_3 < \dots < \lambda_n < \lambda_{n+1} < \dots$

- (a) There is the smallest e' value λ_1
- (b) There is no largest e' value, i.e.

$$\lim_{n \rightarrow \infty} \lambda_n = +\infty$$

3. To every n -value λ_n , there exists an associated n -function $\phi_n(x)$, defined up to a multiplicative constant ($\neq 0$).

$\phi_n(x)$ has exactly $n-1$ roots in $a < x < b$

4. The n -functions $\phi_n(x)$ form a complete set in a sense that any piecewise smooth function $f(x)$ can be represented by a generalized

Fourier series

$$f(x) \sim \sum_{n=1}^{\infty} a_n \phi_n(x)$$

Furthermore, this series converges to

$$\frac{1}{2} [f(x^+) + f(x^-)] \quad \text{for } a < x < b$$

5. The eigenfunctions belonging to different e ' values are orthogonal relative to the weight function $\bar{v}(x)$:

$$\int_a^b \phi_n(x) \phi_m(x) \bar{v}(x) dx = 0 \quad \text{if } n \neq m$$

$\langle \phi_n, \phi_m \rangle$: inner product of ϕ_n and ϕ_m

6. Any eigenvalue can be related to eigenfunction through Rayleigh Quotient:

$$\lambda = \frac{-p\phi\phi' / a^b + \int_a^b [p(\phi')^2 - q\phi^2] dx}{\int_a^b \phi^2 \sigma dx}$$

Consider a special case:

Sturm-Liouville problem: $\phi'' + \lambda\phi = 0$

$$\phi(0) = 0, \quad \phi(L) = 0$$

$p(x) = 1, \quad q(x) = 0, \quad \sigma(x) = 1, \quad a = 0, \quad b = L$
 $\lambda_n = \left(\frac{n\pi}{L}\right)^2 > 0$ and $n = 1, 2, \dots$

1. ... We showed previously that they are real eigenvalues $\lambda_n = \left(\frac{n\pi}{L}\right)^2$

Real eigenvalues: we showed that otherwise there was no nontrivial function $\phi(x)$.

$$f(x) \sim \sum_{n=1}^{\infty} a_n \phi_n(x) \quad / \quad \phi_m(x) \delta(x)$$

$$\int_a^b$$

$$\int_a^b f(x) \phi_m(x) \delta(x) dx = \int_a^b \sum_{n=1}^{\infty} a_n \phi_n(x) \phi_m(x) \delta(x) dx$$

$$= \sum_{n=1}^{\infty} a_n \int_a^b \phi_n(x) \phi_m(x) \delta(x) dx =$$

$= 0$ if $n \neq m$

$$= a_m \int_a^b \phi_m^2(x) \delta(x) dx$$

$$a_m = \frac{\int_a^b f(x) \phi_m(x) \delta(x) dx}{\int_a^b \phi_m^2(x) \delta(x) dx}$$

\therefore

We assumed that λ_n were real but now with Sturm-Liouville Thm we know that λ 's are real (Thm proves that).

2. Ordering e'values:

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2, \quad n=1, 2, \dots$$

$$\lambda_1 < \lambda_2 < \dots < \lambda_n < \lambda_{n+1} < \dots$$

$\lambda_1 = \left(\frac{\pi}{L}\right)^2$: the smallest e'value

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2 \rightarrow \infty \text{ as } n \rightarrow \infty$$

3. Sine functions: $\phi_n(x) = \sin \frac{n\pi x}{L}$, $0 \leq x \leq L$

In the interval $0 < x < L$, $\sin \frac{n\pi x}{L}$ has exactly $n-1$

roots.

4. Complete set: any piecewise smooth function $f(x)$ can be represented as a Fourier series:

$$f(x) \sim \sum_{n=1}^{\infty} a_n \sin \frac{n\pi x}{L} = \sum_{n=1}^{\infty} a_n \phi_n(x)$$

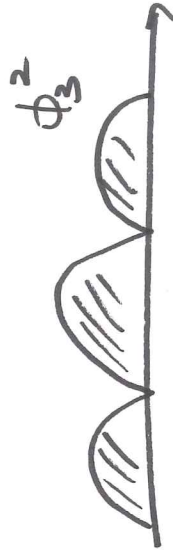
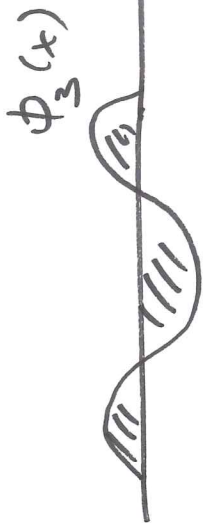
5. Orthogonality of sine functions

In general,

$$\int_a^b \phi_n(x) \phi_m(x) dx = 0 \quad \text{if } n \neq m$$

$m \rightarrow n$

$$a_n = \frac{\int_a^b f(x) \phi_n(x) \sigma(x) dx}{\int_a^b \phi_n^2(x) \sigma(x) dx}$$



Our special case: $\phi_n(x) = \sin \frac{n\pi x}{L}$

$$\sigma(x) = 1$$

$$a=0, \quad b=L$$

$$a_n = \frac{\int_0^L f(x) \sin \frac{n\pi x}{L} \cdot 1 dx}{\int_0^L \sin^2 \frac{n\pi x}{L} \cdot 1 dx}$$

$$\int_0^L \sin^2 \frac{n\pi x}{L} dx = \frac{L}{2}$$

$$\therefore a_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi x}{L} dx : \text{ same as before}$$

6. Rayleigh Quotient We will prove later that e 's value

λ and associated e 's function $\phi(x)$ are related by

$$\lambda = \frac{-p \phi \phi' \Big|_a^b + \int_a^b [p(\phi')^2 - q\phi^2] dx}{\int_a^b \phi^2 dx}$$

$$\lambda = \frac{\int_a^b \phi^2 q dx}{\int_a^b \phi^2 dx}$$

Special case: $\phi(0) = \phi(L) = 0$ $p=1, q=0, \delta=1$

$$\therefore \lambda = \frac{0 + \int_0^L (\phi')^2 dx}{\int_0^L \phi^2 dx} > 0$$

$$(\phi')^2 \geq 0 \Rightarrow \int_0^L (\phi')^2 dx \geq 0 \Rightarrow \boxed{\lambda \geq 0}$$

Q Can λ be 0?

\nearrow Let $\lambda = 0$ be an e 'value

$$0 = \frac{\int_0^L (\phi')^2 dx}{\int_0^L \phi^2 dx} \Rightarrow \int_0^L (\phi')^2 dx \geq 0 \Rightarrow \phi' \equiv 0$$

$\rightarrow \phi = \text{const}$ but $\phi(0) = \phi(L) = 0 \Rightarrow \phi \equiv 0$

$\therefore \lambda = 0$ is not an ϵ -value.

Hence, $\lambda > 0$

Note: we knowed that ϵ -values $\lambda > 0$ without knowing $\phi(x)$ or λ explicitly!