

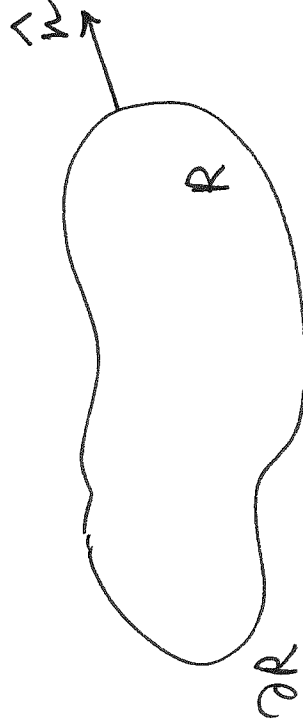
1.5 Heat Equation in Higher Dimensions

Consider a region R in $3D$

R : region

∂R : boundary of R

\hat{n} : outward unit normal to the boundary

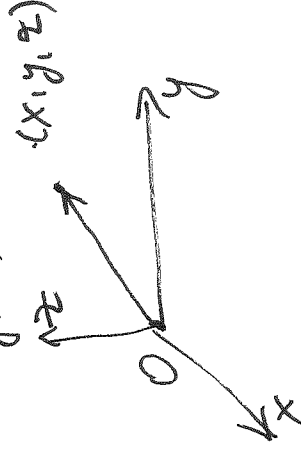


The total thermal energy in R is

$$\iiint_R c(\vec{x}) \rho(\vec{x}) u(\vec{x}, t) dV$$

where $dV = dx dy dz$

$$\vec{x} = (x, y, z)$$



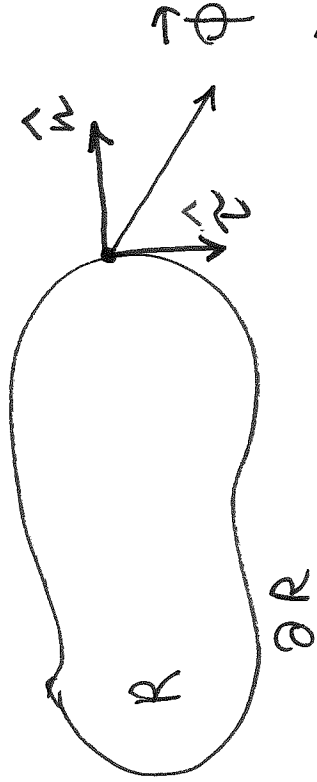
Conservation of thermal energy:

rate of change of total thermal energy = heat energy flowing across the boundary per unit time + energy generated in volume region R due to sources/sinks per unit time

$$\text{rate of change of total thermal energy} = \frac{d}{dt} \iiint_R c(\vec{x}) \rho(\vec{x}) u(\vec{x}, t) dV$$

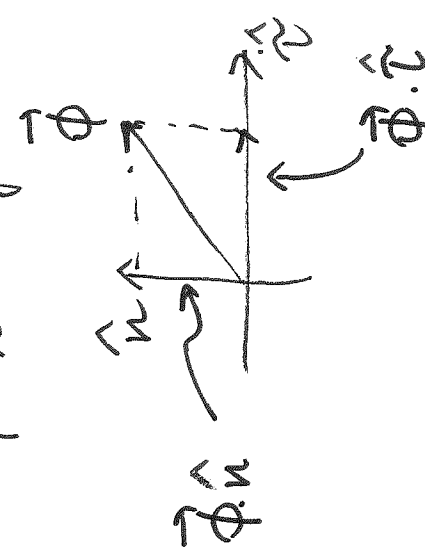
(temperature)

Let $\vec{\Phi}$ be a heat flux vector that represents heat energy that flows per unit time per unit area. The magnitude of $\vec{\Phi}$ is the thermal energy that flows across the boundary.

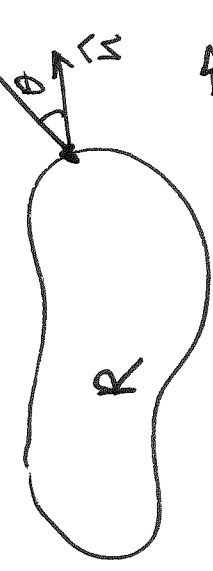


\hat{n} : outward unit normal
 \hat{t} : unit tangent vector

Tangent component of $\vec{\Phi}$, i.e. $\vec{\Phi} \cdot \hat{t}$ has no contribution to the amount of energy that goes through the boundary. For \hat{t} , there is nothing that crosses the boundary. Only normal component $\vec{\Phi} \cdot \hat{n}$ of heat flux vector $\vec{\Phi}$ contributes.



θ : angle between $\vec{\Phi}$ and \hat{n}



Component of $\vec{\Phi}$ in the normal direction \hat{n} is $\vec{\Phi} \cdot \hat{n} = |\vec{\Phi}| \cdot |\hat{n}| \cdot \cos \theta = |\vec{\Phi}| \cdot \cos \theta$

Amount of thermal energy that crosses the boundary ∂R is

$$\oiint_{\partial R} \vec{\Phi} \cdot \vec{n} \, dS$$

∂R

Heat flux is said to be positive if heat flows

in the direction of outward unit normal \vec{n} .

In the conservation of energy equation, if heat flows out the region R , the total energy decreases

$$\Rightarrow \text{amount of energy that flows across the boundary} = - \oiint_{\partial R} \vec{\Phi} \cdot \vec{n} \, dS$$

∂R

The thermal energy due to sources/sinks is

$$\iiint_R Q(\vec{x}) dV$$

Conservation of energy:

$$\frac{d}{dt} \underbrace{\iiint_R c(\vec{x}) \rho(\vec{x}) u(\vec{x}, t) dV}_{\text{"}} = - \iint_{\partial R} \vec{\Phi} \cdot \hat{n} dS + \iiint_R Q(\vec{x}) dV$$

$$\iiint_R c(x) \rho(x) \frac{\partial u(\vec{x}, t)}{\partial t} dV$$

Gauss divergence Thm

Let $\vec{A} = \langle A_1, A_2, A_3 \rangle = A_1 \hat{i} + A_2 \hat{j} + A_3 \hat{k}$ be continuously differentiable. Then

$$\iiint_R \underbrace{\vec{\nabla} \cdot \vec{A}}_{\text{div } \vec{A}} dV = \iint_{\partial R} \vec{A} \cdot \hat{n} dS$$

divergence of \vec{A}

where $\vec{\nabla} = \frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k} = \left\langle \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right\rangle$: gradient vector

$\nabla u = \langle u_x, u_y, u_z \rangle$: gradient of u

$$\text{div } \vec{A} = \vec{\nabla} \cdot \vec{A} = \left\langle \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right\rangle \cdot \langle A_1, A_2, A_3 \rangle = \frac{\partial A_1}{\partial x} + \frac{\partial A_2}{\partial y} + \frac{\partial A_3}{\partial z}$$

Using divergence Thm, we can write

$$\iint_{\partial R} \vec{\Phi} \cdot \hat{n} \, dS = \iiint_R \nabla \cdot \vec{\Phi} \, dV$$

Note: in 1D we had

$$\phi(a) - \phi(b) = - \int_a^b \frac{\partial \phi}{\partial x} \, dx$$

Then the conservation of energy equation becomes

$$\iiint_R c_p \frac{\partial u}{\partial t} \, dV = - \iiint_R \nabla \cdot \vec{\Phi} \, dV + \iiint_R Q \, dV$$

OR
$$\iiint_R \left[c_p \frac{\partial u}{\partial t} + \nabla \cdot \vec{\Phi} - Q \right] dV = 0$$

since R was an arbitrary region in $\mathbb{R}^3 \Rightarrow$ conservation of energy in differential form

$$c\rho \frac{\partial u}{\partial t} + \nabla \cdot \vec{\Phi} - Q = 0$$

$$c\rho \frac{\partial u}{\partial t} + \nabla \cdot \vec{\Phi} = Q$$

(*)

Fourier's Law of Heat Conduction:

$$\vec{\Phi} = -K_0(\vec{x}) \nabla u$$

where $K_0(\vec{x})$ is heat conductivity coefficient.

$$\phi = -K_0(x) \frac{\partial u}{\partial x}$$

Note: in 1D:

Here $\nabla u = \left\langle \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z} \right\rangle = \frac{\partial u}{\partial x} \hat{i} + \frac{\partial u}{\partial y} \hat{j} + \frac{\partial u}{\partial z} \hat{k}$

Now substitute expression for $\vec{\Phi}$ in $(*)$

$$c_p \frac{\partial u}{\partial t} - \nabla \cdot (k_0 \nabla u) = Q$$

the general heat equation

If material has uniform thermal properties, then c, ρ, k_0 are constant. Let set $Q=0$, then we can write heat equation as

where $K = \frac{k_0}{c_p}$: thermal diffusivity coefficient

$$\frac{\partial u}{\partial t} = K \nabla^2 u$$

$$\begin{aligned} \Delta^2 u &= \nabla \cdot (\nabla u) = \left\langle \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right\rangle \cdot \left\langle \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}, \frac{\partial u}{\partial z} \right\rangle = \\ &= \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\partial u}{\partial y} \right) + \frac{\partial}{\partial z} \left(\frac{\partial u}{\partial z} \right) = \\ &= \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \end{aligned}$$

Laplacian of u

To solve the heat equation, we need to specify an initial condition and boundary conditions.

Initial condition: $u(x, y, z, 0) = f(x, y, z)$: initial temperature

where $(x, y, z) \in \mathbb{R}$

Boundary conditions:

1. Prescribed temperature,

$$u(x, y, z, t) = T(x, y, z, t), \quad \text{where } (x, y, z) \in \partial R$$

2. Given flux:

$$-k_0 \nabla u \cdot \hat{n} = \phi(x, y, z, t), \quad (x, y, z) \in \partial R$$

For insulated boundary

$$\nabla u \cdot \hat{n} = 0 \quad \text{on } (x, y, z) \in \partial R$$

Note : $\nabla u \cdot \hat{n} = \frac{\partial u}{\partial n}$: directional derivative of u in direction of unit vector \hat{n}