

Principle of Linear Superposition

Recall, a ^{linear} \forall DE $\mathcal{L}u = f$ is homogeneous if $f=0$, i.e. $\mathcal{L}u=0$.

If u_1 and u_2 satisfy the same linear homogeneous DE, then their linear combination $c_1u_1 + c_2u_2$ is also a solution of the same equation.

Pf Since u_1, u_2 satisfy the same linear homog. DE, we can write

$\mathcal{L}\{u_1\} = 0$ and $\mathcal{L}\{u_2\} = 0$
 We need to show that $\mathcal{L}\{c_1 u_1 + c_2 u_2\} = 0$.

Indeed,
 $\mathcal{L}\{c_1 u_1 + c_2 u_2\}$ is $\mathcal{L}\{c_1 u_1\} + c_2 \mathcal{L}\{u_2\} = 0$
 linear

$\therefore c_1 u_1 + c_2 u_2$ satisfies the same linear homog. DE

Boundary Conditions

The concept of linearity and homogeneity also applies to BCs.

Examples of linear BCs:

$u(0, t) = f(t)$: linear nonhomog. BC

$u_x(L, t) = g(t)$: linear nonhomog. BC

$u_x(0, t) = 0$: linear homog. BC

- $K_0(L) u_x(L, t) = H [u(L, t) - g(t)]$: linear nonhomog. BC

Nonlinear BC:

$$u_x(L, t) = u^2(L, t)$$

↑ nonlinear

Aside

$y = ax + b$: linear function
linear term

$y = ax^2 + bx + c$: quadratic
 x^2

Def a homogeneous (linear) BC is the condition satisfied by a trivial solution ($u \equiv 0$)

$u(x, t)$: nonlinear

$\sqrt{u(x, t)}$: nonlinear
 $\ln u_x$: —

Chapter 2 Separation of Variables

Consider the linear homogeneous 1D heat equation with homogeneous BCs:

$$u_t = k u_{xx} \quad 0 \leq x \leq L$$

$$u(0, t) = 0, \quad u(L, t) = 0$$

$$u(x, 0) = f(x)$$

Note To solve a nonhomogeneous problem, we need to learn first how to solve a homog. problem.

Separation of variables

Look for solutions of the form

$$u(x,t) = \phi(x) G(t)$$

Separation of variables, introduced by Bernoulli, allows one to reduce PDE to ODE.

Note Separation of variables can only be used on linear homogeneous DEs w/ linear homogeneous BCs.

$$\frac{\partial u}{\partial t} = \phi(x) \frac{dG}{dt}$$

$$\frac{\partial^2 u}{\partial x^2} = \frac{d^2 \phi}{dx^2} G(t)$$

Substitute these derivatives in the heat eqⁿ

$$u_t = k u_{xx} :$$

$$\phi(x) \frac{dG}{dt} = k \cdot \frac{d^2\phi}{dx^2} G(t) \quad \left| \quad \frac{1}{k \phi(x) G(t)} \right.$$

$$\frac{dG}{dt} = \frac{d^2\phi}{dx^2} \phi(x) = -\lambda$$

$$\underbrace{\frac{d^2\phi}{dx^2} \phi(x)}_{\substack{\text{f}^n \text{ of } x \\ \text{alone}}} = \underbrace{-\lambda}_{\substack{\text{separation constant}}} \quad (\text{separation constant})$$

i.e. both sides should be equal to some constant, called a separation constant.

We will "ignore" IC for now.

We obtain two ODEs:

$$\frac{dG}{dt} = -\lambda \quad \text{or}$$

$$k \frac{dG}{dt} + \lambda k G(t) = 0$$

time dependent problem for $G(t)$

$$\frac{d^2\phi}{dx^2} = -\lambda \quad \text{or}$$

$$\frac{d^2\phi}{dx^2} + \lambda \phi = 0$$

x dependent problem for $\phi(x)$

The goal is to construct a nontrivial solution, i.e. nonzero solution.

Boundary conditions

$$u(0, t) = 0 \Rightarrow \phi(0) G(t) = 0 \Rightarrow$$

$$\boxed{\phi(0) = 0}$$

$\neq 0$ since we need a nontrivial solution

$$\boxed{\phi(L) = 0}$$

$$u(L, t) = 0 \Rightarrow \phi(L) \underbrace{G(t)}_{\neq 0} = 0 \Rightarrow$$

The time-dependent problem for $G(t)$:

$$\frac{dG}{dt} + aK G(t) = 0 \quad \text{separable ODE}$$

$$G = 0 \text{ is also a sol}^n$$

$$\frac{dG}{dt} = -aK G(t)$$

$$\frac{dG}{G} = -aK dt \quad G \neq 0$$

$$\int \frac{dG}{G} = -aK \int dt$$

$$\ln |G| = -aKt + \tilde{C} \quad | \text{ exp } \quad e^{-aKt} = X$$

$$|G| = e^{-akt} + \tilde{C} = e^{-akt} \cdot e^{\tilde{C}} > 0$$

to combine w/ $G \geq 0$ we can write

$$G(t) = C e^{-akt}$$

C can be > 0 , < 0 and $= 0$

C : arbitrary const
Boundary Value Problem for $\phi(x)$ (BVP)

$$\frac{d^2\phi}{dx^2} + \lambda\phi(x) = 0$$

$$\phi(0) = 0, \quad \phi(L) = 0$$

This is called an eigenvalue problem.

λ is an eigenvalue and $\phi(x) \neq 0$ is an associated eigenfunction.

Def λ is called an eigenvalue if there exists a nontrivial function $\phi(x)$, called an eigenfunction, that satisfies the above BVP.

We need to find all possible λ 's for which we can find nontrivial $\phi(x)$

Case I λ 'values and e' functions for $\lambda > 0$

$$\frac{d^2\phi}{dx^2} + \lambda\phi(x) = 0, \quad \lambda > 0$$

$$x'' + 4x = 0$$

$$x = e^{rx}$$

$$(D^2 + 4)x = 0$$

Assume $\phi(x) = e^{rx}$. $\Rightarrow \phi' = r e^{rx}$, $\phi'' = r^2 e^{rx}$

$$r^2 e^{rx} + \lambda \cdot e^{rx} = 0$$

$e^{rx} (r^2 + \lambda) = 0 \Rightarrow r^2 + \lambda = 0$: characteristic eq^{1/2}
 $\neq 0$ $r^2 = -\lambda$ $\lambda > 0$

$$r = \pm i\sqrt{\lambda}$$

\Rightarrow solutions are $e^{i\sqrt{\lambda}x}$, $e^{-i\sqrt{\lambda}x}$ \equiv

$$\cos\sqrt{\lambda}x + i\sin\sqrt{\lambda}x, \quad \cos\sqrt{\lambda}x - i\sin\sqrt{\lambda}x$$

In fact, real & imaginary parts, $\cos\sqrt{\lambda}x$, $\sin\sqrt{\lambda}x$

are also solutions \Rightarrow

operator approach:

$$(D^2 + \lambda)\phi = 0$$

$$\pm i\sqrt{\lambda}$$

$$\phi(x) = C_1 \cos\sqrt{\lambda}x + C_2 \sin\sqrt{\lambda}x$$

$$\phi = C_1 \cos\sqrt{\lambda}x + C_2 \sin\sqrt{\lambda}x$$

$$\phi(0) = 0 \Rightarrow c_1 \underbrace{\cos 0}_1 + c_2 \cancel{\sin 0} = 0 \Rightarrow \boxed{c_1 = 0}$$

$$\phi(x) = c_2 \sin \sqrt{\lambda} x$$

$$\phi(L) = 0 \Rightarrow c_2 \sin(\sqrt{\lambda} L) = 0 \Rightarrow \sin(\sqrt{\lambda} L) = 0$$

#0 for a nontrivial solution

$$\sqrt{\lambda} L = n\pi, \quad n = 1, 2, 3, \dots$$

$$\Rightarrow \lambda_n = \left(\frac{n\pi}{L}\right)^2, \quad n = 1, 2, 3, \dots \quad \text{eigenvalues}$$

$$\phi_n(x) = \sin \sqrt{\lambda_n} x = \sin \frac{n\pi x}{L} \quad ; \quad \text{eigenfunctions}$$