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We showed last time that under scaling

$$F(x, t) = \alpha F(\alpha x, \alpha^2 t)$$

the ratios $\frac{(\alpha x)^2}{\alpha^2 t} = \frac{x^2}{t}$ or $\frac{\alpha x}{\sqrt{\alpha^2 t}} = \frac{x}{\sqrt{t}}$

and $\frac{\alpha F}{\alpha x} = \frac{F}{x}$ or $\frac{\alpha F}{\sqrt{\alpha^2 t}} = \frac{F}{\sqrt{t}}$

stay the same, hence $F(x, t)$ can be written as

$$F(x, t) = \frac{1}{\sqrt{t}} f\left(\frac{x}{\sqrt{t}}\right) \quad \text{or} \quad F(x, t) = \frac{1}{\sqrt{t}} g\left(\frac{x^2}{t}\right)$$

or $F(x, t) = \frac{1}{x} h\left(\frac{x}{\sqrt{t}}\right)$

The non-uniqueness is not a problem, and we can choose any of these functions f , g or h . For example, let

$$F(x, t) = \frac{1}{\sqrt{t}} f\left(\frac{x}{\sqrt{t}}\right), \quad \frac{x}{\sqrt{t}} = \xi$$

$F(x, t)$ satisfies

$$\frac{F}{t} - F_{xx} = \delta(x) \delta(t)$$

If $t \neq 0 \Rightarrow F_t - F_{xx} = 0$ (1) homog. PDE

$$F_t = \frac{f'(\xi) \frac{\partial \xi}{\partial t} \cdot \sqrt{t} - f(\xi) \cdot \frac{1}{2\sqrt{t}}}{t} \quad (\equiv)$$

$$\xi = \frac{x}{\sqrt{t}}$$

$$\xi = \frac{x}{\sqrt{t}}$$

$$\frac{\partial \xi}{\partial t} = -\frac{1}{2} \frac{x}{t^{3/2}}$$

$$\frac{\partial \xi}{\partial x} = \frac{1}{\sqrt{t}}$$

$$\equiv \frac{f'(\xi) \cdot \left(-\frac{1}{2}\right) \frac{x}{t \cdot \sqrt{t}} \sqrt{t} - f(\xi) \frac{1}{2\sqrt{t}}}{t} =$$

$$= \frac{1}{t} \left[-\frac{x f'}{2t} - \frac{f}{2t^{3/2}} \right] = -\frac{x f'}{2t^2} - \frac{f}{2t^{3/2}}$$

$$F_x = \frac{1}{\sqrt{t}} f'(\xi) \cdot \underbrace{\frac{\partial \xi}{\partial x}}_{\frac{1}{\sqrt{t}}} = \frac{1}{t} f'(\xi)$$

$$F_{xx} = \frac{1}{t^{3/2}} f''(\xi)$$

Substitute F_t, F_{xx} into (1).

$$\underbrace{\left(-t^{3/2}\right) \cdot \left[-\frac{x f'}{2t^2} - \frac{f}{2t^{3/2}} \right]}_{F_t} - \underbrace{\left(-t^{3/2}\right) \cdot \left[\frac{1}{t^{3/2}} f''(\xi) \right]}_{F_{xx}} = 0 \quad \text{if } t \neq 0$$

$$f'' + \frac{x}{2\sqrt{x}} f' + \frac{f}{2} = 0 \quad t \neq 0 \quad \zeta = \frac{x}{\sqrt{x}}$$

$$(2) \quad f'' + \frac{\zeta}{2} f' + \frac{f}{2} = 0 \quad \text{ODE for } f \text{ 2nd order}$$

Note

$$\zeta f' + f = (\zeta f)'$$

⇒ we can write (2) as

$$f'' + \frac{1}{2} (\zeta f)' = 0$$

Integrate once

$$f' + \frac{1}{2} \zeta f = A \quad \begin{array}{l} Q(\zeta) \\ \text{const of integration} \end{array}$$

$$p(\zeta) = e^{\int \frac{1}{2} \zeta d\zeta} = e^{\frac{\zeta^2}{4}}$$

$$p \cdot f = \int p A d\zeta + B \quad \text{const}$$

$$e^{\frac{\zeta^2}{4}} \cdot f = \int e^{\frac{\zeta^2}{4}} \cdot A d\zeta + B \quad \left| e^{-\frac{\zeta^2}{4}} \right.$$

$$f(\zeta) = A e^{-\frac{\zeta^2}{4}} \int e^{\frac{s^2}{4}} ds + B e^{-\frac{\zeta^2}{4}}$$

Aside
Method of integrating factor for 1st order linear ODEs:

$$y' + P(x)y = Q(x)$$

$$p(x) = e^{\int P(x) dx}$$

integrating factor

$$p y = \int p Q dx + C$$

Solve for y

$$F = \frac{f(\zeta)}{\sqrt{t}}$$

We can find constants A and B by considering the total heat

$$\text{total heat } H(t) = \int_{-\infty}^{\infty} F(x,t) dx \quad \equiv \quad \text{integral of temperature over the entire domain}$$

$$\text{Let } f_1(\zeta) = e^{-\frac{\zeta^2}{4}} \int e^{\frac{s^2}{4}} ds = \left. \begin{array}{l} \sigma = \frac{s^2}{4} \\ s = 2\sqrt{\sigma} \\ ds = \frac{d\sigma}{\sqrt{\sigma}} \end{array} \right| =$$

$$= e^{-\frac{\zeta^2}{4}} \int e^{\sigma} \sigma^{-1/2} d\sigma \quad \begin{array}{l} \text{integration} \\ \text{by parts} \\ \text{twice} \end{array}$$

We can show that

$$(*) \quad f_1(\zeta) = \frac{2}{|\zeta|} + O\left(\frac{1}{\zeta^3}\right) \quad \text{as } |\zeta| \rightarrow \infty$$

$$\zeta = \frac{x}{\sqrt{t}} \quad \equiv \quad \underbrace{\frac{A}{\sqrt{t}} \int_{-\infty}^{\infty} f_1\left(\frac{x}{\sqrt{t}}\right) dx}_{\text{I}} + \frac{B}{\sqrt{t}} \int_{-\infty}^{\infty} e^{-\frac{x^2}{4t}} dx$$

$$f_1(\xi) = f_1\left(\frac{x}{\sqrt{t}}\right) \stackrel{(*)}{=} \frac{2\sqrt{t}}{x} \left(1 + O\left(\frac{1}{\xi^2}\right)\right)$$

not defined
at $x=0$

$\Rightarrow \int_{-\infty}^{\infty} f_1(\xi) d\xi$ blows up at the origine

since it behaves as $\int \frac{dx}{x}$

$$\Rightarrow A=0$$

Hence,

$$H(t) = \frac{B}{\sqrt{t}} \int_{-\infty}^{\infty} e^{-\frac{x^2}{4t}} dx \Rightarrow 1 = \frac{B}{\sqrt{t}} \int_{-\infty}^{\infty} e^{-\frac{x^2}{4t}} dx$$

$$H(t) = \begin{cases} 0, & t < 0 \\ 1, & t > 0 \end{cases}$$

Heaviside function

will justify this
later

Then we can write

$$1 = \frac{B}{\sqrt{t}} \int_{-\infty}^{\infty} e^{-\frac{x^2}{4t}} dx = 2B \int_{-\infty}^{\infty} \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4t}} dx =$$

$$= \left| \xi = \frac{x}{2t} \right| = 2B \underbrace{\int_{-\infty}^{\infty} e^{-\xi^2} d\xi}_{\sqrt{\pi}} = 2B\sqrt{\pi}$$

$$\Rightarrow B = \frac{1}{2\sqrt{\pi}}$$

$$\therefore F(x,t) = B e^{-\frac{x^2}{4t}} \cdot \frac{1}{\sqrt{t}} = \frac{1}{2\sqrt{\pi}} \cdot \frac{1}{\sqrt{t}} e^{-\frac{x^2}{4t}}$$

$$\text{or } \boxed{F(x,t) = \frac{1}{2\sqrt{\pi t}} e^{-\frac{x^2}{4t}}}$$

Show that $H(t)$ above is the Heaviside function.

Recall that the fundamental solution satisfies

$$u_t - u_{xx} = \delta(x)\delta(t)$$

$$H(t) = \int_{-\infty}^{\infty} F(x,t) dx$$

$$F_t - F_{xx} = \delta(x)\delta(t)$$

$$\frac{dH}{dt} = \int_{-\infty}^{\infty} F_t(x,t) dx = \int_{-\infty}^{\infty} [F_{xx}(x,t) + \delta(x)\delta(t)] dx$$

$$= F_x(x, t) \Big|_{-\infty}^{\infty} + \delta(t) \underbrace{\int_{-\infty}^{\infty} \delta(x) dx}_{=1} = F_x(\overset{\nearrow 0}{+\infty}, t) - F_x(\overset{\nearrow 0}{-\infty}, t)$$

since

$u \rightarrow 0$ as $|x| \rightarrow \infty$

u is smooth

$$+ \delta(t) = \delta(t)$$

$\therefore \frac{dH}{dt} = \delta(t) \Rightarrow H(t)$ is a ^{unit} step function

$$\text{i.e. } H(t) = \begin{cases} 0, & t < 0 \\ 1, & t > 0 \end{cases}$$

$$\text{i.e. } H(t) = 1, \quad t > 0.$$